

Series 5000

stereo control preamplifier

Designed as the perfect partner to our Series 5000 MOSFET stereo amplifier, the stereo control preamplifier has been designed with equal attention to quality and detail. This month David Tilbrook explains the principles behind the design of the moving coil and moving magnet stages and gives details on constructing them.

David Tilbrook

JUST AS a loudspeaker represents a non-linear load to the output stage of a power amplifier, a moving magnet or moving coil cartridge represents a non-linear source impedance to the input stage of a preamplifier. This is the cause of many of the problems associated with any preamp.

Both moving coil and moving magnet cartridges generate electrical signals through the interaction of a coil of wire and a magnetic field. The signal voltage produced is therefore proportional to the relative velocity between the coil and the magnet assemblies. This relationship is predicted by Faraday's law of induction, expressed mathematically as:

$$\epsilon = -\frac{d\phi}{dt}$$

where ϵ is the signal voltage at any instant and ϕ is the magnetic flux.

The signal voltage produced at any instant is proportional to the rate of change of flux with respect to time,

$$\text{i.e. } \frac{d\phi}{dt}$$

The design of the cartridge must ensure that a linear relationship exists between the position of the stylus cantilever assembly and the magnetic flux.

Our adaptation of a famous scene from Irving Stone's book about Michelangelo, 'The Agony and the Ecstasy'. During the time Michelangelo was painting his masterpiece frescoes in the Sistine Chapel the Pope continually asked "... when will you make an end?". Likewise, Roger Harrison has kept asking David Tilbrook when the Series 5000 would finish. The answer was the same in both cases! (Thanks to the cartoonist, Brendan Akhurst, and the choirboys: Jack O'Donnell of Altronics, Dick Smith, and Gary Johnston of Jaycar. Collyn Rivers looks on from the sidelines).

In this way changes in the position of the stylus give rise to changes in the magnetic field intensity. So the rate of change of stylus position with respect to time will be proportional to the signal voltage, i.e.:

$$\epsilon \propto \frac{dx}{dt}$$

where ϵ is the signal voltage and x is the stylus displacement from its equilibrium position.

This means that the waveform actually 'on' the grooves is not proportional to the signal voltage itself. Instead it is proportional to the integral of the signal waveform. If a square wave, for example, is to be produced from a record, the waveform as seen in the groove with a microscope will be a triangle wave.

Since the signal voltage is proportional to the velocity of the stylus, the signal slope is proportional to the acceleration of the stylus. In order for high signal slopes to be reproduced accurately by the cartridge it is important that the mass of the stylus cantilever assembly be kept as small as possible. At the same time, however, it is important to realise that the cartridge cantilever assembly and its associated suspension and magnet/coil system form a resonant mass-spring system analogous to a complex electrical series resonant circuit.

At one particular frequency, called the resonant frequency, the impedance of the cartridge will no longer be related linearly to the driving force on the stylus, and distortion results. To over-

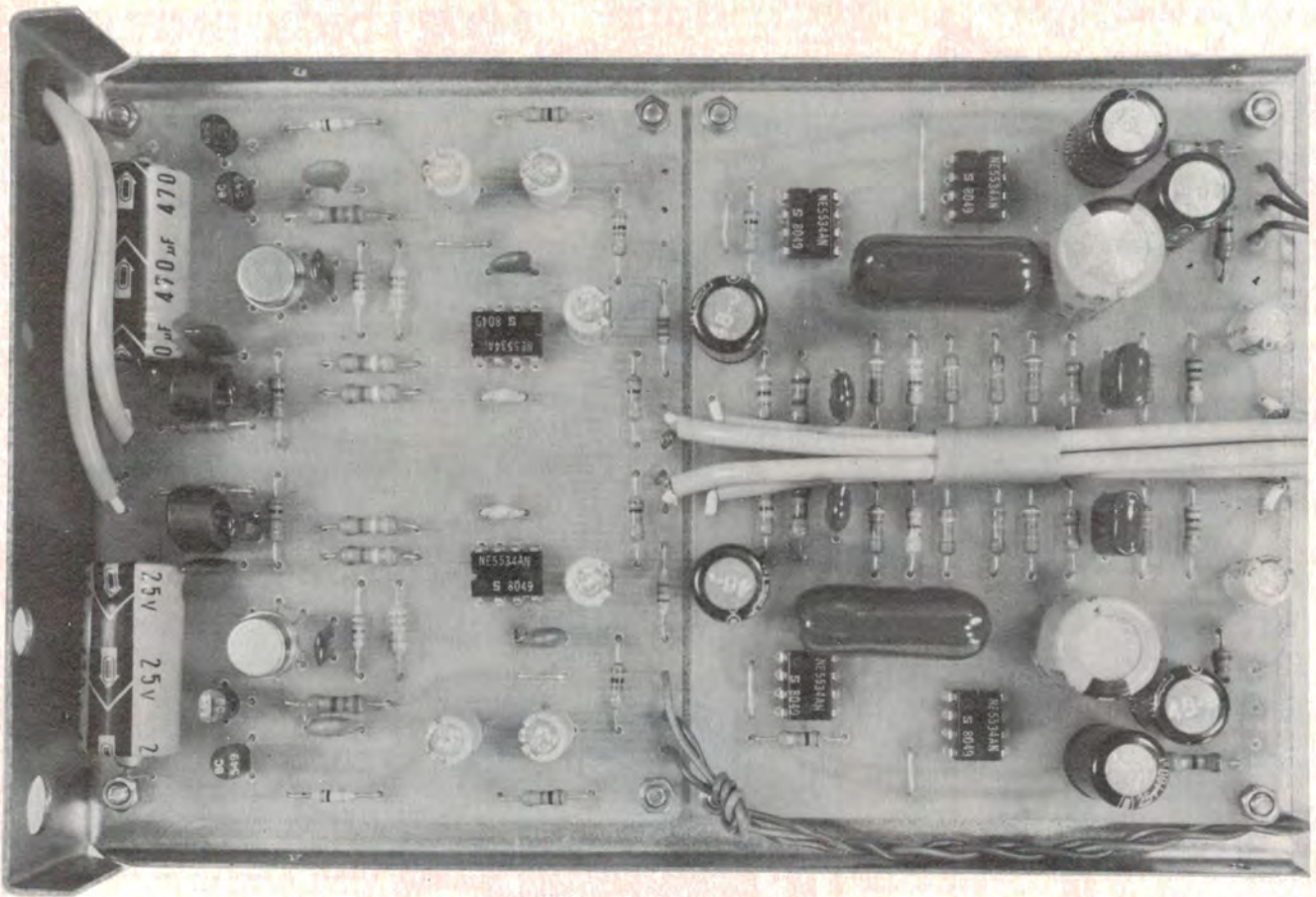
come this problem two techniques are used simultaneously.

Firstly the resonant frequency of the cartridge is moved to a frequency below the audio spectrum. Using the damped mass-spring model of a magnetic cartridge we can predict that the resonant frequency will depend on the mass of the stylus cantilever assembly and on the 'springiness' of the cantilever's suspension. This springiness is characterised by a number, often given the symbol k , called the spring constant. Spring constant is defined in terms of the force needed to bring about a certain compression or extension of the spring. Stiffer springs have a higher value for k . The spring constants associated with magnet cartridges, however, are so small that the numbers are hard to interpret. For this reason cartridge manufacturers usually specify this quantity by quoting the reciprocal of the spring constant, $\frac{1}{k}$, called *compliance*. Stiffer suspension systems have lower compliance figures.

As stated earlier, the cartridge resonant frequency is a function of both the mass and the compliance of the cantilever and suspension system. The damped resonance mass-spring model of a magnetic cartridge predicts that the resonant frequency will be given by the equation:

$$f = \frac{1}{2\pi\sqrt{mC}}$$

where m is the mass of the cantilever/stylus system and C is the compliance of the stylus suspension system.



The low level Series 5000 Preamp input stages — for moving magnet and moving coil cartridges — are housed in a steel enclosure inside the main preamp case. This is a view inside the enclosure; moving coil stage to the left, moving magnet stage at right.

(Notice that the equation for the resonant frequency of magnetic cartridges has exactly the same form as the equation for the resonant frequency of an electrical resonance circuit, i.e:

$$f = \frac{1}{2\pi\sqrt{LC}}$$

where C in this case is capacitance and L is inductance. The fact that a mechanical system like the stylus cantilever assembly of a cartridge should be described by such an equation is another striking example of the consistency of nature.)

The equation predicts that the resonant frequency of the cartridge can be decreased by increasing either the mass or the compliance. Since the mass of the moving parts in the cartridge must be kept small so the stylus can respond quickly to changes in the record groove, the compliance must be increased until a suitably low resonant frequency is obtained. Most high-quality magnetic cartridges have resonant frequencies below 10 Hz.

The second technique used to overcome problems associated with this resonance characteristic is to decrease

the Q of the system by damping the resonance with a suitable combination of mechanical and electrical losses. Mechanical damping is obtained by deliberately introduced friction within the cantilever suspension system. The cantilever is often terminated into a rubber mounting block for this purpose. The electrical damping comes about as a direct consequence of the law of conservation of energy. The cartridge is acting as a generator, delivering power to the input resistance of the pre-amplifier. Since energy is absorbed by this load resistance the Q of the cartridge resonance is decreased.

Until recently most stereo magnetic cartridges consisted of two fixed coils between the poles of a small magnet attached to the cantilever. Modulation of the record groove produces movement of the magnet, changing the magnetic flux and generating the signal voltage.

The coils usually have a large number of turns so that a reasonable signal voltage can be produced (typically in the order of 20 mV). The resistance of these coils usually ranges between 200-1000 ohms, but their impedance can be much higher, especially at high

frequencies where the inductance of the coils becomes important. This type of cartridge is sometimes called a moving magnet cartridge to distinguish it from the more recently developed moving coil types. The relatively high reactive component of the cartridge impedance combined with the effects of the natural cartridge resonances makes it essential that the input impedance of the moving magnet (MM) input stage have well-defined characteristics if best performance is to be obtained from this type of cartridge. Most MM cartridges require a load impedance consisting of 47k of resistance shunted by several hundred picofarads. This capacitance is often provided by the shielded cable, but most cartridges require some additional capacitance across the MM input. In exceptional cases the input capacitance due to the shielded cable is too high. In these cases the length of the shielded cable used between the turntable and the MM amp should be decreased until the correct value is achieved.

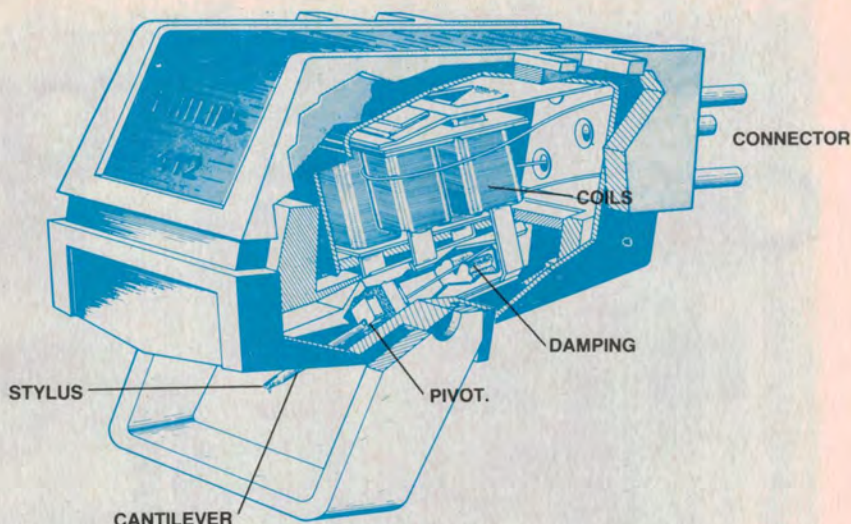
In order to obtain the flattest frequency response possible from an MM cartridge it is essential that the

stereo control preamp

load resistance be constant over the complete audio spectrum *and beyond*. For this reason measurements done on the input resistance of MM amps at one particular frequency (usually 1 kHz) are practically useless. Many input stages exhibit a characteristic of falling input resistance at high frequencies. The input resistance of a bipolar transistor, for example, even with a small amount of emitter current, is insufficient to ensure a constant resistive load to an MM cartridge. The common two or three transistor phono stages of a few years ago often suffered badly from this problem, degrading the top end performance of an otherwise good MM cartridge. The problem occurs because all bipolar transistors have decreasing gain at high frequencies.

The most common method used to increase the input impedance of a bipolar input stage is through the use of negative feedback. The decrease in gain of the individual transistors in the stage at high frequencies decreases the overall open loop gain of the stage, which in turn decreases the amount of negative feedback available. Furthermore, the negative feedback is often applied at the emitter of the first transistor. The problem with this con-

View inside a typical high-quality moving magnet cartridge, the Philips GP412.



figuration is that the phase response in the negative feedback loop can easily be affected by the complex reactances of the cartridge and connecting cables, producing unwanted frequency response variations, or even instability in some cases.

All these problems come under the

general heading of 'cartridge impedance interaction', and represent the most important single reason for the difference in sound between preamplifiers. Most preamps suffer from some degree of cartridge impedance interaction and in many cases the effects are pronounced. ▶

SPECIFICATIONS

ETI-478MM MOVING MAGNET INPUT STAGE

Gain: 74, 1 kHz

Frequency response: Conforms to RIAA Equalisation ± 0.2 dB. (This is the performance of the prototype. The actual figure obtained will be determined by the accuracy and long-term stability of the components used.)

Total harmonic distortion: $< 0.001\%$, 1 kHz, 10 mV RMS input

Headroom: > 28 dB with respect to 5 mV RMS input signal, i.e. 135 mV RMS max.

Noise: Total equivalent input noise, 122 nV 'A', input shorted, 216 nV flat, input shorted.

S/N ratio:

	1 mV	5 mV	10 mV
Flat	73 dB	87 dB	93 dB
A-weighted	78 dB	92 dB	98 dB

IDEAL RIAA

Hz	dB
2	-0.2
4	+5.7
8	+11.2
16	+15.4
20	+16.3
30	+17.0
40	+16.8
50	+16.3
80	+14.2
100	+12.9
150	+10.3
200	+8.2
300	+5.5
400	+3.8
500	+2.6
800	+0.7
1k	0.0
1k5	-1.4
2k	-2.6
3k	-4.8
4k	-6.6
5k	-8.2
6k	-9.6
8k	-11.9
10k	-13.7
15k	-17.2
20k	-19.6

MEASURED - SERIES 5000

dB
-0.2
+5.7
+11.2
+15.4
+16.2
+17.0
+16.8
+16.2
+14.2
+12.8
+10.2
+8.1
+5.4
+3.7
+2.6
+0.7
0.0
-1.3
-2.4
-4.7
-6.6
-8.1
-9.6
-11.9
-13.8
-17.1
-19.5

SPECIFICATIONS

ETI-478MC MOVING COIL INPUT STAGE

Gain: 24

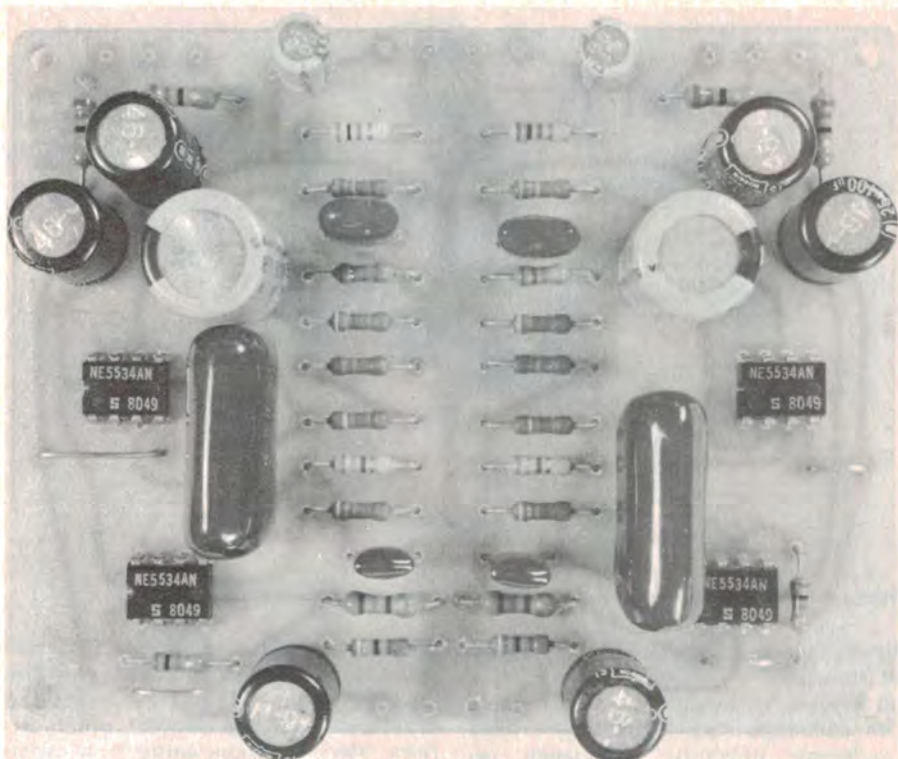
Frequency response: 7 Hz-135 kHz $+0, -1$ dB

Total harmonic distortion: $< 0.003\%$, 1 kHz, 30 mV input

Noise: Total equivalent input noise: 83 nV flat, input shorted. 42 nV 'A', input shorted. 56 nV flat, after RIAA Eq., input shorted. 34 nV 'A', after RIAA Eq., input shorted.

S/N ratio of MC input stage after RIAA Equalisation:

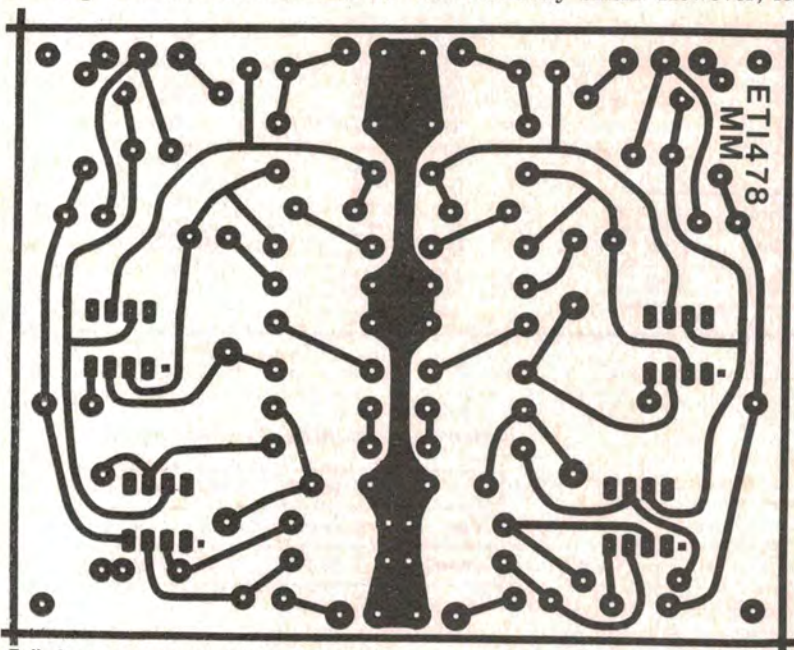
	60 μ V	200 μ V	500 μ V
Flat	61 dB	71 dB	79 dB
A-weighted	65 dB	75 dB	83 dB



The completed moving magnet stereo input stage (ETI-478MM). Note that this project may also be incorporated in existing equipment if you wish.

The Series 5000 stereo control preamp has been designed specifically to overcome the problem of cartridge impedance interaction. This has been achieved by separating the MM input stage into two separate active stages (see Figure 1). The first stage consists of a single NE5534AN configured as a linear amplifier with a closed loop gain of around 8.3. The large amount of overall negative feedback increases the

input impedance of the stage so that the measured input impedance is simply that of the 470k resistor, R2. Since the 5534 has a small signal bandwidth of around 10 MHz without additional compensation, the input impedance will remain unchanged over a very wide frequency range. The high input impedance of this stage would usually allow the input capacitor C2 to be conveniently small. However, for best



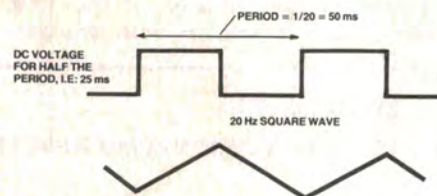
Full-size pc board artwork for the moving magnet stage.

noise performance the value must be increased substantially. This is covered in detail later in this article.

Capacitor C2 is necessary since it is not advisable to allow dc current from the first stage to flow through the cartridge. The value of C2 used here is 100μ , and this sets the lower -3 dB point well below 1 Hz. The upper -3 dB point of this stage is well above 100 kHz. An extended frequency response is necessary so that the accuracy of the RIAA equalisation is not affected by frequency response variations that might otherwise occur in the first stage.

RIAA equalisation

We said earlier that the signal voltage produced by a magnetic cartridge is proportional to the velocity of the stylus. If a low frequency signal is to be reproduced by a magnetic cartridge, large excursions of the stylus are necessary. If for example a 20 Hz square wave is to be reproduced by the cartridge then the cartridge must produce a dc voltage at its output for a period of 25 ms.



In order to do this the stylus must move at a constant speed for this period of time, and therefore the waveform in the record groove is a triangle wave, as stated earlier.

Typical output voltages from moving magnet cartridges are in the order of 1 mV-2 mV for a stylus velocity of 1 cm/sec. So if the peak voltage required on the square wave was, say, 10 mV, a stylus velocity of 10 cm/sec would be required for a medium-sensitivity cartridge, so the stylus must move at a constant speed of 10 cm/sec for a 25 ms time interval. The stylus therefore moves a total distance of 2.5 mm! On a stereo record the channels are cut in opposite walls of the record groove. If a low frequency mono signal is to be produced, both sides of the record groove force the stylus away from its equilibrium position, and a large vertical stylus excursion results. In the case of our square wave, the vertical excursion would be roughly 3.5 mm, which is simply not possible. The record would have to be as thick as most turntable platters!

Two measures are used to overcome

stereo control preamp

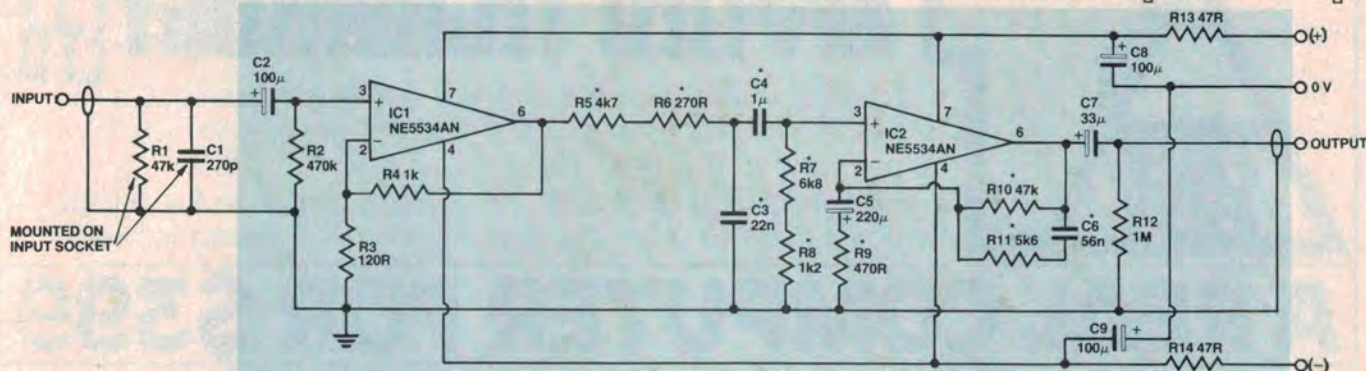


Figure 1. Circuit of one channel of the moving magnet input stage (ETI-478MM). Note that the RIAA equalisation is incorporated in this stage. Components for the other channel are designated R101, C101, IC101, etc.

* DENOTES COMPONENTS ASSOCIATED WITH THE RIAA EQUALISATION

this problem. Firstly the two channels are recorded on the record 180° out of phase, so that the large vertical excursion is replaced by a large horizontal excursion. Secondly, the low end of the frequency response is attenuated before the recording process, so the stylus excursions are decreased. The specific amount of low frequency attenuation is defined as that which would be caused by a first-order high-pass filter with a time constant of 318 µs (i.e: the filter would be formed by an ideal resistor/capacitor filter, in which $R \times C = 318 \mu s$). To convert from these time constants into frequency, simply apply the equation:

$$f = \frac{1}{2\pi t} \quad (t = \text{time constant})$$

This is equivalent to a 6 dB/octave filter with a -3 dB point at 500 Hz. To prevent the low end from rolling off indefinitely a second 6 dB/octave filter is used to flatten the response again at 3150 µs or 50 Hz. After this equalisation is applied, the stylus excursion of the 20 Hz square wave, for example, is decreased to around 0.3 mm, which is manageable.

Similar problems occur at very high frequencies. If we consider now a 20 kHz square wave at the same output voltage and hence the same recording velocity, the stylus now only moves a total distance of 2.5 µm! Such minute distances are only a few orders of magnitude larger than the surface irregularities in the vinyl, so at these frequencies the signal to noise ratio is poor. To overcome this problem the top end is recorded at a higher level, which increases the stylus excursions and thereby improves the signal to noise ratio. The modifications to the recorded frequency response are referred to as RIAA pre-emphasis or equalisation (RIAA stands for Recording Institute Association of America), and must be corrected for by the input stage. The

RIAA playback equalisation must boost the bass end and attenuate the treble end of the audio spectrum to return the overall frequency response to that of a linear system.

Since the low end is amplified most of all by the RIAA playback signal, any turntable rumble or cartridge/turntable resonances will be amplified. Modern power amps are quite capable of delivering full power to a pair of loudspeakers at 10 Hz or below, so appreciable amounts of subsonic content can be fed to the loudspeaker. This is potentially dangerous to the bass driver and decreases the clarity and accuracy of the low end.

In an attempt to overcome this problem the RIAA has proposed a change to its playback equalisation curve. The extreme bass frequencies are attenuated on playback by the addition of another time constant. This takes the form of a single-pole RC filter with a time constant of 7950 µs, i.e: a -3 dB point of 20 Hz. Since the frequency response is already flattened by the 3150 µs time constant, this new time constant gives a 6 dB attenuation rate below about 20 Hz. The resulting RIAA playback equalisation is shown in Figure 2. Note that there are four time constants associated with the proposed RIAA equalisation: 7950 µs, 3150 µs,

HOW IT WORKS

ETI-468MM

The input from a moving magnet cartridge is connected to the non-inverting input of an NE5534AN via capacitor C2. R2 provides a dc current path to the input of the differential pair in the op-amp. The gain of this stage is determined by the ratio R4 to R3, which is around 8.3 in this case.

The resistor R1 provides a fixed resistive load necessary for best performance from an MM cartridge. Most cartridge manufacturers recommend that the input resistance be shunted by a certain amount of capacitance. This is the purpose of capacitor C1, the value of which should suit most cartridges. If you wish to optimise the value of this capacitor don't forget to allow several hundred picofarads for the shielded cable capacitance.

The best way to ensure that the cartridge is loaded correctly is with a test record containing a square wave track, and an oscilloscope. With the correct cartridge load and a good tonearm/cartridge combination, a good square wave can be obtained.

The value of resistor R1 at 47k is effectively in parallel with R2, giving an input resistance of 43k, slightly below the 47k normally used for MM input stages. This is unimportant, however, and will not affect performance of the cartridge. The important thing is that the value of this resistance remain constant over the full audio spectrum and beyond. In any case the value of the input resistance is easily changed by increasing the value of R1 to, say, 56k instead of 47k.

The output of the first stage is fed to two

6 dB/octave RC filters that provide one half of the RIAA equalisation. Resistors R5, R6 and capacitor C3 form a first-order low-pass filter set at the 75 µs time constant of the RIAA curve. At these frequencies (around 2122 Hz) the 1 µF capacitor appears as a short circuit connecting R7 and R8 in parallel with the capacitor C3. This must be compensated when choosing the value of C3 to ensure the correct RIAA equalisation. Similarly C4, R7 and R8 form a low frequency high-pass filter set at 20 Hz (the 7950 µs time constant).

The output of these two filters is fed to the input of the second op-amp stage. The remaining RIAA equalisation is accomplished by the feedback loop around this stage. At frequencies below 500 Hz the 56nF capacitor C6 has a relatively high impedance. The voltage gain is therefore determined by resistors R9 and R10. At higher frequencies, however, where the impedance of C6 is less, both resistors R10 and R11 are in circuit. The capacitor C5 decreases the gain, at dc, of the second stage to unity, ensuring a low dc offset at the output and therefore symmetrical output stage clipping.

The 1M resistor R12 ensures that the dc voltage on the output remains at 0 V. This is important so that operation of the selector switch following the stage will not cause thumps in the output.

Resistors R13, R14 and capacitors C8, C9 isolate the supply to the stage to decrease the effects of interactions between stages and to ensure freedom from 50 Hz ripple.

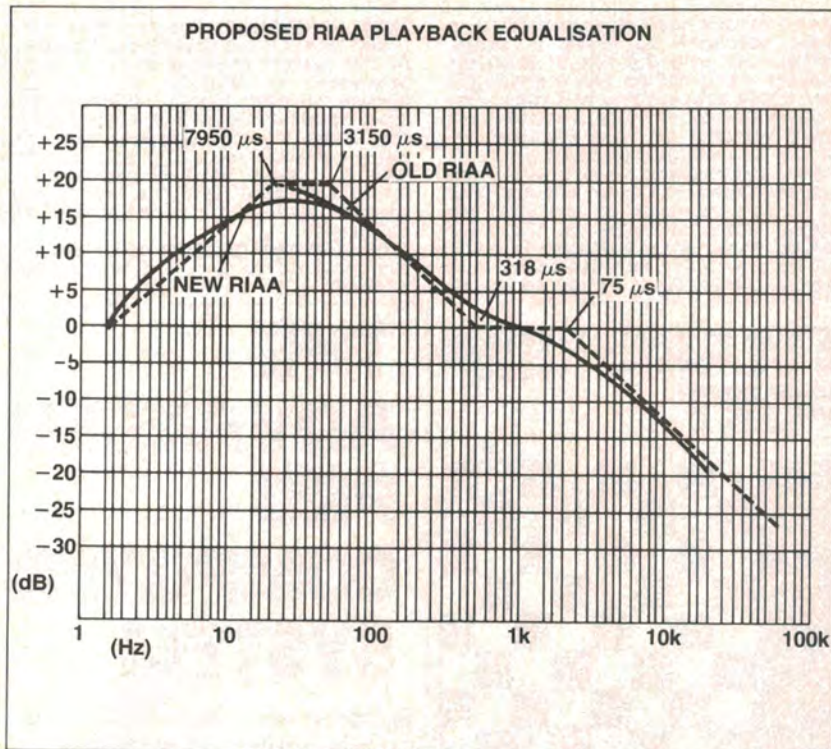
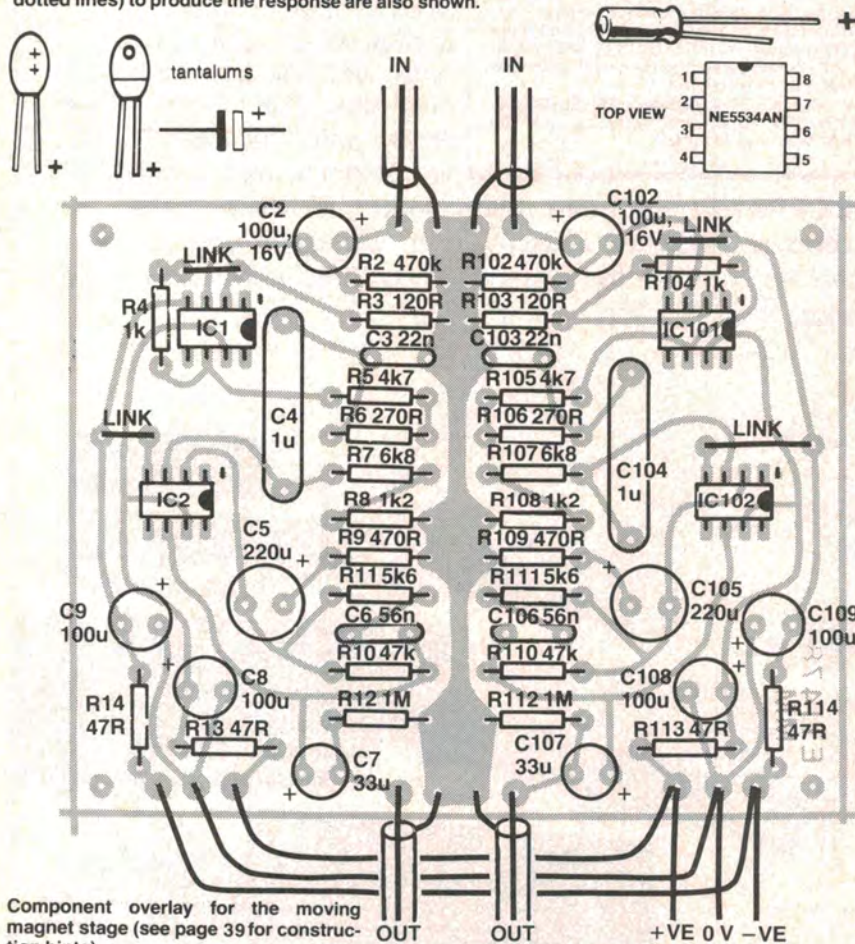


Figure 2. Old and 'new' RIAA equalisation curves (solid line). The individual time constants (Bode plot — dotted lines) to produce the response are also shown.



Component overlay for the moving magnet stage (see page 39 for construction hints).

318 μ s and 75 μ s. These are shown on the Bode plot, which is the dotted line in Figure 2. It should be emphasised, however, that the introduction of this low frequency time constant is not sufficient to remove severe cases of turntable or tonearm resonance. Some preamps incorporate multiple-order subsonic filters that offer a very fast roll-off below 20 Hz. The problem with this, however, is that severe cases of tonearm resonance or rumble generate distortion harmonics well above 20 Hz, into the audio spectrum. The only real cure is to remove the problem at the turntable or tonearm.

Many different techniques are used to give the preamp the desired equalisation. The most common is to include the RIAA equalisation circuitry into the feedback loop of the first stage. Figure 3 shows a very simple MM input stage of the general type often found in medium-priced amplifiers.

Transistor Q1 functions as a standard common emitter amplifier offering a voltage gain that is determined by the total impedance from its collector to earth divided by the total impedance from its emitter to earth. Transistor Q2 is a PNP transistor but functions in an identical manner. The product of their two voltage gains is called the open loop gain of the stage. If a current path is now made available from the output of Q2 back to the emitter of Q1, the voltage

PARTS LIST — ETI-478MM FOR STEREO PC BOARD

Resistors

R1, R101	all 1/4W metal film
R1, R101	47k
R2, R102	470k
R3, R103	120R
R4, R104	1k
R5, R105	4k7, 1%
R6, R106	270R, 1%
R7, R107	6k8, 1%
R8, R108	1k2, 1%
R9, R109	470R, 1%
R10, R110	47k, 1%
R11, R111	5k6, 1%
R12, R112	1M
R13, R113, R14, R114	47R

Capacitors

C1, C101	270p ceramic
C2, C102	100 μ , 16 V electro.
C3, C103	22n greencap
C4, C104	1 μ greencap
C5, C105	220 μ , 16 V electro.
C6, C106	56n greencap
C7, C107	33 μ , 25 V electro.
C8, C108, C9, C109	100 μ , 25 V electro.

Integrated circuits

IC1, IC101	NE5534AN
IC2, IC102	NE5534AN

Miscellaneous

1 x ETI-478MM pc board; assorted mounting hardware; shielded cable.

stereo control preamp

gain will now drop to a new figure called the closed loop gain. This is negative feedback, and it has the effect of decreasing the distortion and increasing the input impedance of the stage. (See Series 5000 MOSFET Power Amp articles in ETI Jan., Feb., March 1981 for more information on negative feedback.)

The RIAA equalisation is introduced by applying the negative feedback via a network with a frequency dependent impedance. However, since this stage relies on the presence of negative feedback to ensure a satisfactorily high input impedance, the input impedance will vary as a function of frequency. The cartridge, however, must be loaded by a constant resistance if cartridge impedance interaction is to be avoided. Furthermore, since the negative feedback is coupled to the complex output impedance of the cartridge via the base-emitter junction of Q1, the negative feedback and hence the frequency response of the stage can be affected by the cartridge itself. As a result this type of stage can suffer badly from cartridge impedance interaction.

In the development of the Series 5000 preamp several input stage configurations were tested for noise, distortion and cartridge impedance interaction. When a medium-priced moving magnet cartridge was connected to a stage like that in Figure 3, severe cartridge impedance interaction was evident. The frequency response of the preamplifier peaked above 2 dB at 13 kHz and rolled off rapidly above 15 kHz. The same cartridge when connected to the Series 5000 MM amp exhibited quite a good frequency response to beyond 20 kHz, and the frequency response graph obtained was identical to that when a FET buffer amp was placed between the cartridge and the input stage, indicating almost total lack of cartridge impedance interaction in the 5000 stage. This is a result of the use of the separate linear gain stage formed by IC1 (see Figure 1) to isolate the cartridge from the RIAA equalisation.

The Series 5000 Preamp conforms to the proposed RIAA equalisation in Figure 2. The 75 μ s and 7950 μ s time constants are obtained by passive RC filters at the output of the first stage. Resistors R5, R6 and capacitor C3 form a simple 6 dB/octave low-pass filter with a -3 dB point at 2122 Hz, and

$$t = \frac{1}{2\pi f} = \frac{1}{2\pi(2122)} \doteq 75 \mu\text{s}.$$

Capacitor C4, together with resistors R7 and R8, form a 6 dB/octave high-pass filter with a -3 dB point at 20 Hz,

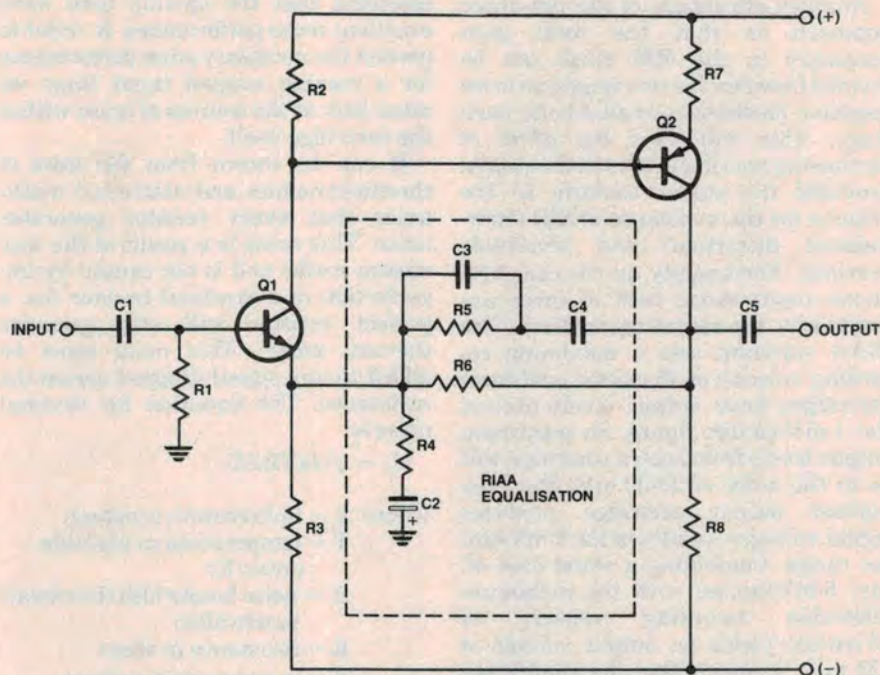


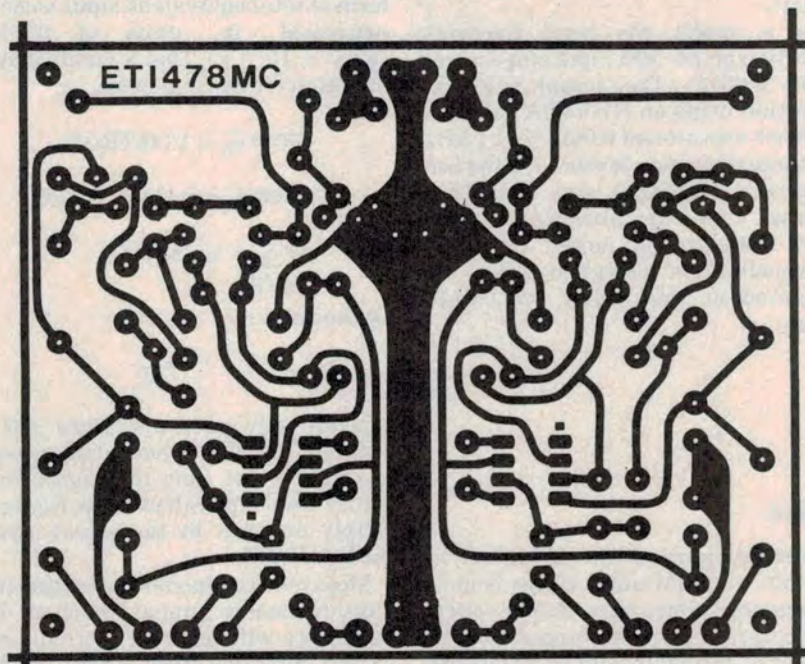
Figure 3. Typical moving magnet input stage found in medium-priced amplifiers.

which is equivalent to a 7950 μ s time constant. The two remaining time constants are introduced into the negative feedback of IC2 and are formed by the values of resistors R9, R10, R11 and capacitor C6.

This method of generating the RIAA curve offers a number of advantages over the more conventional method.

Firstly there is low interaction between the different time constants, so that the RIAA curve can be optimised for a particular cartridge more easily by

changing the resistor or capacitor values slightly. If the 75 μ s time constant is included in the negative feedback of a stage, the gain of the stage must decrease to unity at a suitably high frequency, so the stage must be compensated for unity gain to prevent instability. In the Series 5000 MM stage the gain of the second stage does not drop below 10; since the NE5534AN is internally compensated for gains of 3 or above, no additional compensation is required.



Full-size artwork for the moving coil input stage.

Another advantage of the two-stage approach is that the total gain necessary in the MM stage can be divided between the two stages, so more negative feedback is available for each stage. This will have the effect of decreasing non-linearities in the stages, provided the stages conform to the criteria for the avoidance of SID (slew-induced distortion) and amplitude overload. Fortunately, in the case of a phono input stage, both of these are limited by the recording medium. The RIAA standard sets a maximum recording velocity of 25 cm/sec, and most cartridges have output levels around the 1 mV/cm/sec figure. So maximum output levels from such a cartridge will be in the order of 20-30 mV. Even the highest output cartridge produces signal voltages usually in the 5 mV/cm/sec range. Combining a worst case of, say, 5 mV/cm/sec with the maximum allowable recording velocity of 25 cm/sec yields an output voltage of 125 mV. To ensure that the input stage cannot be overloaded we simply set the gain of these stages so that this maximum input signal cannot drive the output of the input stages into clipping. The NE5534AN is capable of driving to within 2 V of the supply voltage, so a supply voltage of ± 15 V gives the desired gain of around 75. We have divided this gain between the two input stages so that the first stage has a gain of 8.3 and the second stage a gain of 9 in the midband region (the actual gain of the second stage is of course a function of frequency due to the RIAA equalisation).

As a result the total harmonic distortion of this MM input stage is well under 0.001%. The actual measured distortion using an HP3580A spectrum analyser was around 0.0005% at 1 kHz. (At these distortion levels even the best distortion analysers are practically useless, since the distortion is well below the level of noise.) Similarly, intermodulation distortion (IMD) was measured at well below the 0.001% figure.

Noise

Another very important parameter for both MC and MM input stages is noise performance. Since an op-amp is used as the first stage of the MM input amp, we have only limited control over the noise performance of the stage. It is therefore

essential that the op-amp used have excellent noise performance. In order to predict the necessary noise performance for a moving magnet input stage we must look at the sources of noise within the cartridge itself.

It can be shown from the laws of thermodynamics and statistical mechanics that every resistor generates noise. This noise is a result of the way nature works and is not caused by imperfection in a practical resistor (i.e. a perfect resistor will still generate thermal noise). This noise must be added to any signal dropped across the resistance. The equation for thermal noise is:

$$\bar{e}_n = \sqrt{(4kTR\Delta f)}$$

where k = Boltzmann's constant,
 T = temperature in absolute units (K)
 Δf = noise bandwidth (brickwall bandwidth)
 R = resistance in ohms
 \bar{e}_n = average noise voltage

This equation predicts that thermal noise is raised by increasing resistance temperature or the bandwidth of the measuring equipment. So the frequency response of the apparatus used to determine thermal noise must be quoted if the figure is to be meaningful. Furthermore, the Δf here refers to a 'brickwall frequency response', not the usual half-power bandwidth, although for many purposes this is sufficiently accurate. To overcome this problem noise performance is often quoted in the form of total equivalent input noise and expressed in units of nV/ $\sqrt{\text{Hz}}$ (1 nV = 10^{-9} V). This is justified by the equation for thermal noise, i.e.:

$$\text{since } \bar{e}_n = \sqrt{(4kTR\Delta f)}$$

$$\text{then } \bar{e}_n = (\sqrt{\Delta f})(\sqrt{(4kTR)})$$

$$\text{or } \frac{\bar{e}_n}{\sqrt{\Delta f}} = \sqrt{(4kTR)}$$

So the ratio:

$$\frac{\bar{e}_n}{\sqrt{\Delta f}}$$

depends only on temperature and resistance, and this is just what we want. In order to get from this figure to an actual total equivalent noise figure we simply multiply by the square root of the bandwidth.

Most moving magnet cartridges have a coil resistance around 500 ohms. This resistance will generate thermal noise, so the cartridge itself limits the best possible signal-to-noise ratio. Using the

equation for thermal noise we obtain for the noise of the cartridge:

$$\frac{\bar{e}_n}{\sqrt{\text{Hz}}} = \sqrt{(4 \times 1.37 \times 10^{-23} \times 290 \times 500)}$$

(assuming temperature of resistor is $\approx 290\text{K}$).

$$\text{i.e. } \frac{\bar{e}_n}{\sqrt{\text{Hz}}} = 2.8 \times 10^{-9}$$

$$\text{i.e. } \bar{e}_n = 2.8 \text{ nV}/\sqrt{\text{Hz}}$$

We can express this in more familiar terms by converting the cartridge noise figures into a signal-to-noise ratio figure. In audio we can regard the bandwidth in question to be around 20 kHz, i.e. $\sqrt{\Delta f} = 140$, and $140 \times 2.8 \text{ nV}/\sqrt{\text{Hz}} = 392 \text{ nV}$. If the average output level of the cartridge is around 5 mV, the signal-to-noise ratio is given by:

$$20 \log \frac{5 \times 10^{-3}}{392 \times 10^{-9}} \approx 82 \text{ dB}$$

This figure represents the best signal-to-noise ratio possible with most moving magnet cartridges, since this is due to noise generated within the cartridge itself. A well-designed input stage should approach this noise figure as closely as possible without sacrificing performance in other equally important parameters such as distortion and frequency response.

The noise generated by an active device is determined by a number of factors, the most important of which is the current flowing through the device. However, since we have elected to use a high-quality operational amplifier for the input stage, we have no control over device current. All we can do is choose a low-noise op-amp and avoid degrading its noise figure as much as possible. The NE5534AN has a recommended equivalent input noise voltage around 4 nV/ $\sqrt{\text{Hz}}$, only 3 dB above the noise generated by the cartridge itself! In order not to degrade this figure we must keep all resistances in series with the cartridge as low as possible. Any additional resistance will generate a thermal noise voltage of its own, which must be added vectorally to that generated by the cartridge. From the basic equation of thermal noise generated by two individual resistors R_1 and R_2 , for example, we obtain:

$$\frac{\bar{e}_{n1}}{\sqrt{\Delta f}} = \sqrt{(4kTR_1\Delta f)}$$

$$\frac{\bar{e}_{n2}}{\sqrt{\Delta f}} = \sqrt{(4kTR_2\Delta f)}$$

Here we assume that both resistances are at the same temperature. Since these noise voltages are not correlated (i.e. they consist of 'randomly' changing voltage) we add them using the vector sum:

$$\text{i.e. } \overline{e_{nT}^2} = \overline{e_{n1}^2} + \overline{e_{n2}^2}$$

where $\overline{e_{nT}^2}$ is the square of the total equivalent noise voltage.

$$\text{Therefore } \overline{e_{nT}^2} = 4kT\Delta f(R_1 + R_2)$$

$$\text{or } \overline{e_{nT}} = \sqrt{(4kT\Delta f(R_1 + R_2))}$$

If R_1 now represents the cartridge resistance and R_2 the value of an added resistance equal to the value of R_1 , we get:

$$\overline{e_{nT}} = \sqrt{(4kT\Delta f(2R_1))} = \sqrt{2}\sqrt{(4kT\Delta fR_1)}$$

$$\text{or } \overline{e_{nT}} = 1.4\overline{e_{nT}}$$

equivalent to a 3 dB decrease in the signal-to-noise ratio.

Figure 4 shows the standard technique for connecting an op-amp to a signal generator such as a moving magnet cartridge. Most op-amps, and certainly the 5534, have input stages that consist of a differential pair, providing both inverting and non-inverting inputs.

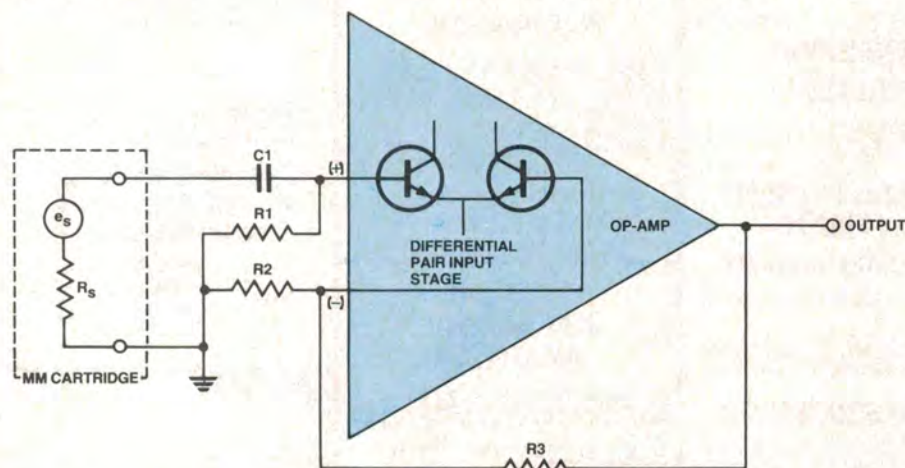


Figure 4. Standard technique for connecting an op-amp to a signal source.

The effective signal voltage generator of the cartridge is represented by e_s and the cartridge resistance by R_s . Resistor R_1 in this case would be 47k, so that the cartridge would have the correct load resistance. (The input impedance of the op-amp is very high and can be ignored for this discussion.) Capacitor C_1 prevents any dc current flowing through the cartridge from the non-inverting input. Since the combination of R_1 and C_1 forms a 6 dB/octave high-pass filter, the value of C_1 would ordinarily be chosen so that the resulting -3 dB point was well below the audio spectrum, around 5 Hz for example. This will occur when the impedance of C_1 is equal to that of R_1 , i.e.:

47k. Since the reactance of the capacitor is given by the equation:

$$X_c = \frac{1}{2\pi fC}$$

we have:

$$C = \frac{1}{2\pi fX_c}$$

$$\text{In this case } C = \frac{1}{2\pi \times 5 \times 47 \times 10^3}$$

$$\doteq 6.77 \times 10^{-7} \text{ Farads.}$$

So to obtain an adequately flat frequency response a suitable value for C_1 would be 680 nF (0.68μF), which is convenient since a greencap could be used.

When noise considerations are taken into account, however, this value is entirely unsuitable. The increasing impedance of C_1 at low frequencies, while not sufficient to cause gross frequency response errors, will seriously degrade the noise performance of the stage. At sufficiently low frequencies the impedance seen by the non-inverting input will be simply the value of R_1 . Using the

So a value around 100 μF should suffice. Notice that this capacitor would have to be an electrolytic or tantalum. Tantalum capacitors are not recommended, however, since their capacitance can be modulated by the input signal, producing considerable distortion at low frequencies.

The value of resistor R_2 must also be low, so that the source impedance to the inverting input of the op-amp can be kept as low as possible. The limitation here is due to the minimum load impedance allowable on the output of the op-amp. Since the gain of the stage is given by the equation:

$$A_v = \frac{R_2 + R_3}{R_2}$$

the ratio of R_2 and R_3 is determined by the desired voltage gain. At the same time, however, the total resistance $R_2 + R_3$ represents the load on the output stage of the op-amp. Since this must not be less than a certain specified resistance, determined by the individual op-amp used, a minimum value for R_2 is predicated. In the Series 5000 MM input stage, for example, the required voltage gain in the first stage is around 8.3, so:

$$\frac{R_2 + R_3}{R_2} = 8.3$$

The NE5534AN has a measured minimum load impedance of 600 ohms, and for minimum distortion it is desirable to increase this slightly, for example to around 1k2. Therefore:

$$\frac{1k2}{R_2} = 8.3 \text{ or } R_2 \doteq 144R$$

A suitable value for R_2 would be 120 ohms, making R_3 1k to give the required voltage gain. Fortunately this value for R_2 is low enough not to have significant effect on the noise performance.

Similar measures must be adopted around the second stage. At low frequencies the non-inverting input of IC2 (see Figure 1) has an input source resistance determined by R_7 and R_8 , i.e. around 8k. The noise performance of the second stage would be improved if this value could be decreased. Unfortunately this would entail increasing the value of C_4 , which is not practical since this capacitor must be a greencap if the preamp is to conform accurately to the RIAA curve. This is not really a problem, however, since the voltage gain in the first stage increases the signal voltage at the input of IC2 to around 40 mV for a 5 mV input signal, ensuring a sufficiently good signal-to-noise ratio in the second stage. ▶

equation for thermal noise given earlier, we can calculate the resulting signal-to-noise ratio. Since

$$\overline{e_n} = \sqrt{(4kT\Delta fR_1)} \\ \overline{e_n} = \sqrt{(4 \times 1.37 \times 10^{-23} \times 290 \times 20 \times 10^3 \times 47 \times 10^3)} = 3.87 \mu V,$$

only 62 dB below 5 mV.

Furthermore, since the input stage is a noise generator, a low source impedance is necessary to minimise the resulting noise at the output of the op-amp. To overcome this problem we increase the value of C_1 so that at worst its impedance at, say, 3 Hz is comparable to that of the cartridge, i.e.:

$$C = \frac{1}{2\pi \times 3 \times 500} = 106 \times 10^{-6} F.$$

Series 5000

The moving coil input stage

The subject of noise performance is particularly important for a moving coil input stage. The moving coil cartridge works on exactly the same principle as the moving magnet. The signal voltages produced are the result of relative motion between a coil of wire and a magnetic flux. In this case, however, the magnet assembly is mounted rigidly to the cartridge body and the coils are mounted on the cantilever assembly; hence the name 'moving coil'.

In order for the total mass and therefore the inertia of the stylus/cantilever system to be kept to a minimum, the coils are made with very fine wire and a small number of turns. Typical output voltages for moving coil cartridges vary widely from one manufacturer to another, but a figure of $40 \mu\text{V/cm/sec}$ is probably a reasonable compromise. A gain of 25 is therefore required to boost this voltage to that of a typical moving magnet cartridge. Once again we can calculate the best possible signal-to-noise ratio for a moving coil cartridge based on its thermal noise. The coil resistance of a moving coil cartridge with an output of $40 \mu\text{V/cm/sec}$ would be approximately 20 ohms (although this figure can vary widely, typically 5-50 ohms).

From the equation for thermal noise we obtain:

$$\frac{\bar{e}_n}{\sqrt{\text{Hz}}} = \sqrt{(4kTR)},$$

$$\text{i.e. } \frac{\bar{e}_n}{\sqrt{\text{Hz}}} = \sqrt{(4 \times 1.37 \times 10^{-23} \times 290 \times 20)}$$

$$\doteq 0.56 \text{ nV}/\sqrt{\text{Hz}}.$$

The total noise over a 20 kHz noise bandwidth is therefore:

$$0.56 \text{ nV} \times \sqrt{(20 \times 10^3)}$$

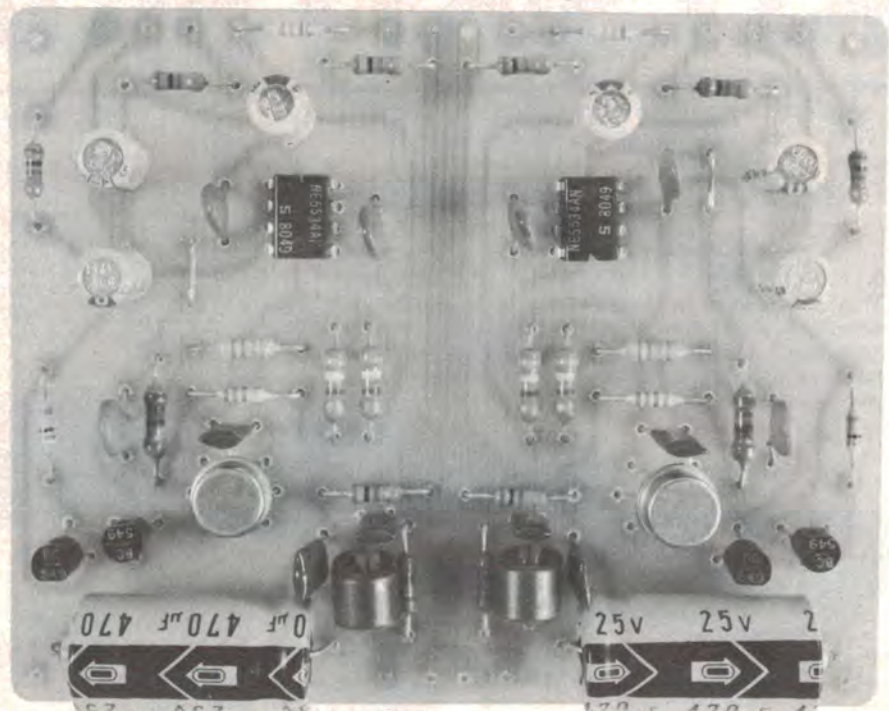
$$\text{i.e. } 0.56 \text{ nV} \times 140 \doteq 78 \text{ nV}.$$

Since the cartridge output voltage will be around $40 \mu\text{V/cm/sec} \times 5 \text{ cm/sec}$, i.e. $200 \mu\text{V}$ for a recording velocity of 5 cm/sec, the resulting signal-to-noise ratio will be:

$$20 \log \frac{200 \times 10^{-6}}{78 \times 10^{-9}}$$

or around 68 dB unweighted.

This figure is only approximate, of course, but it is roughly correct and represents the best possible signal-to-noise ratio with a moving coil cartridge. The object is to design a preamplifier that will approach this noise figure and maintain a flat frequency response, low distortion and a constant resistive input impedance. At these noise levels we cannot use an NE5534AN in a circuit like the MM input stage. The total



The completed moving coil input stage (ETI-478MC). Note that this project could be used 'stand-alone' as an MC head amp.

equivalent input noise in that case was around $4 \text{ nV}/\sqrt{\text{Hz}}$, i.e. 560 nV. The resulting signal-to-noise ratio would be only 51 dB with respect to an input signal of $200 \mu\text{V}$.

In order to achieve a satisfactory noise performance it is necessary to look at the various sources of noise in bipolar transistors and decrease the total equivalent input noise through optimum biasing of the input stage and choice of the first transistor.

One source of noise in the transistor is of course thermal noise. We saw before that to minimise thermal noise it was necessary to ensure a low source resistance over as broad a frequency range as possible. In order to do this for the MC stage the total resistance in series with the source must be kept to a similar value to the source resistance, i.e. around 10 or 20 ohms, depending on the cartridge.

The problem is that the resistance of the base-emitter junction of most bipolar transistors, called the base spreading resistance, is usually much higher than this. One solution is to use a large number of low-noise transistors in parallel to form the input transistor, thus decreasing the base spreading resistance. This was the technique used in the Series 4000 moving coil preamp, published in ETI October 1979. Another solution is to use a power transistor, such as a 2N3055, as the input transistor, and the results using this method can be quite good! The third alternative, and the one we elected to use in this

design, is to make use of an exceptional matched pair produced by National Semiconductor. This device, the LM394, has a low base resistance, very low noise and a high h_{FE} of around 500. (A data sheet for the LM394 will be included in next month's article).

Another source of noise in bipolar transistors is shot noise or base current noise. This is a white noise generator (i.e. the average amplitude of the noise current is constant with frequency), but the noise is increased if emitter current is increased. The base resistance, however, is also a function of the current flowing in the emitter, and is given roughly by the equation:

$$r_b = \frac{26}{I_{E(\text{mA})}}$$

The resistance of the base decreases with increasing emitter current, so noise voltage produced by thermal noise across the base resistance is decreased by increasing the emitter current. In a bipolar transistor, therefore, we have two distinct sources of noise, one increasing with emitter current while the other decreases. For this reason an optimum emitter current exists which represents the best compromise between these two noise sources. With an LM394 operated from source resistances typical of moving coil cartridges, the optimum emitter current is around 8 mA, much higher than would normally be used in an input stage. The result, however, is a very low

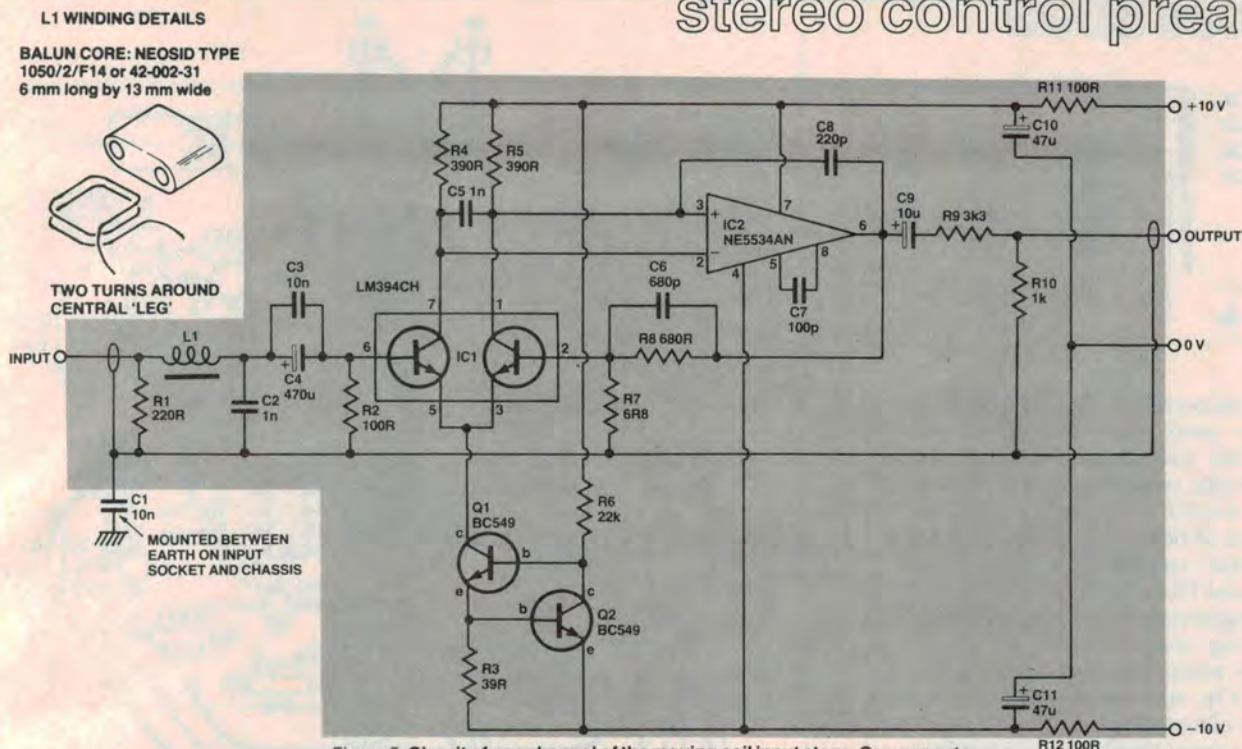


Figure 5. Circuit of one channel of the moving coil input stage. Components for the other channel are designated R101, C101, IC101, etc. Note that data for the LM394 and NE5534 devices are included at the end of this article.

value of input noise for source resistances around 10 ohms.

The complete circuit diagram for the moving coil input stage is shown in Figure 5. The collectors of the LM394 are connected to the input of an NE5534, which functions as a high-gain differential amplifier, providing adequate open loop gain to ensure low distortion and a flat frequency response when negative feedback is applied. The input choke is used to minimise the stage's susceptibility to RF noise.

The input impedance of the stage is determined by the parallel combination of R1 and R2, around 65 ohms for the values shown. This should be suitable for most moving coil cartridges, but is easily changed if required. The dc operating point of the LM394 is determined by the constant current source formed by Q1, Q2, R3 and R6. So the current in resistor R2 is determined by this constant current source and the dc current gain of the LM394. Hence the value of R2 can be increased, in order to increase the input impedance, over a fairly wide range of values without affecting the operation of the circuit.

Once again the input coupling capacitor C4 is used to prevent dc current from flowing through the cartridge. Capacitor C4 is shunted by C3, a 10n capacitor, so that the base of the first transistor in the LM394 is decoupled for RF, through C2. Capacitor C2 represents a shunt capacitance to ensure correct loading of the moving coil cartridge. The value shown should

be suitable for most cartridges, but can be changed for optimisation with any particular cartridge.

To prevent loading the 5534A, the feedback resistor R8 is kept above 600R, i.e.: 680R. Resistor R7 effectively increases with the cartridge and must be kept as low as possible for best noise performance. The value of 6R8 chosen

gives the stage a gain of around 100, which is too high. This is corrected, however, by a simple passive voltage divider at the output, formed by R9 and R10. Capacitor C9 doubles as a feedback isolation capacitor to ensure that reactive components in the load cannot cause a phase shift sufficient to cause oscillation.

HOW IT WORKS ETI-478MC

The input from a moving coil cartridge is fed via L1 and capacitors C3 and C4 to the base of one of the transistors in the LM394, which functions as a differential input stage.

Q1 and Q2 form a constant current source, which stabilises the dc operating point and ensures a high impedance source to the emitters of the differential pair. The constant current source works by ensuring that a constant voltage is maintained across a fixed value of resistance. Resistor R3 is used for this purpose, with the base emitter voltage of Q2 expressed across it. If the current through R3 were to try to increase even slightly, the voltage on the base of Q2 would be increased, turning Q2 on harder. This causes the voltage on the collector of Q2 to decrease, decreasing the current through R3. So Q2 provides negative feedback acting to correct any deviations in the current flowing through the differential pair.

The collectors of the LM394 are shunted by the 1n capacitor C5. This decreases the gain of the first stage at high frequencies and helps to ensure stability (i.e.: freedom from high frequency oscillations).

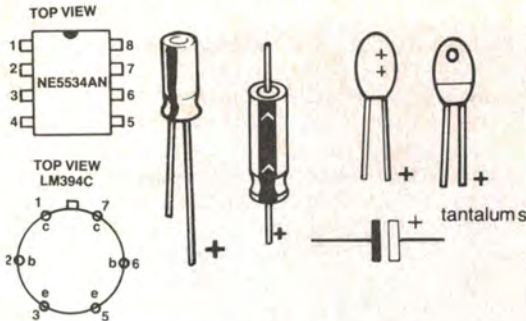
The input stage is operated in full differential mode by connecting both collectors to inputs of the NE5534AN. If this is not done the voltage gain of the input stage is decreased and the signal-to-noise ratio is degraded. Because differential pairs have two base-emitter junctions

in the input circuit, their total equivalent input noise is inferior to that of a single transistor. However, since it is possible using a differential pair to obtain noise figures of the same order of magnitude as the thermal noise of the cartridge, the marginal decrease in the theoretically best signal-to-noise ratio is of little consequence. On the other hand the inherent linearity of a differential pair offers a significant advantage over a single transistor, improving both distortion and high frequency stability.

Capacitor C7 ensures stability of the op-amp by providing adequate compensation for the increased gain around the stage due to the differential pair. C9 provides dc isolation of the stage. The resistors R9 and R10 form a potential divider to decrease the signal level to that suitable for the MM input. If the particular moving coil cartridge used requires a different amount of voltage gain than is provided, the value of R9 can be changed accordingly. Replacing R9 with a short circuit (i.e.: a piece of tinned copper wire in place of the resistor on the circuit board) increases the voltage gain of the stage to slightly over 100.

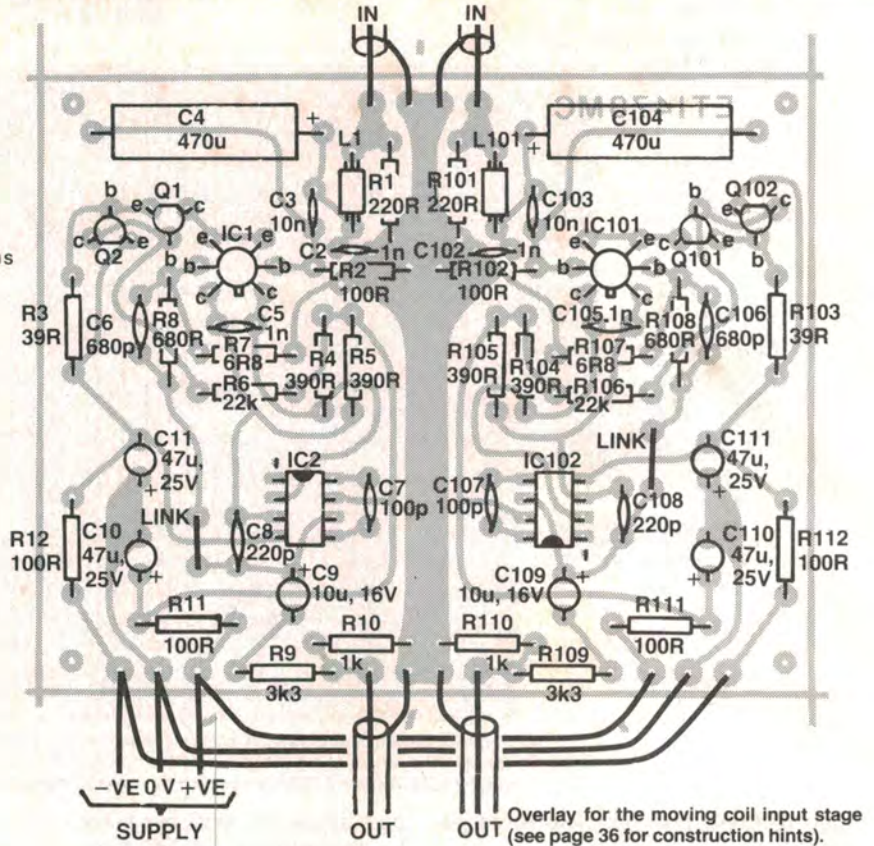
The two RC networks R11, C10 and R12, C11 provide isolation of the supply voltage from other stages using the same power supply. This decreases interactions between stages, thereby improving crosstalk and the overall stability of the preamplifier.

Series 5000



The noise performance of the stage is extremely good. The total equivalent input noise was measured at 83 nV over a 20 kHz noise bandwidth. This is equivalent to 0.6 nV/√Hz or a signal-to-noise ratio of 68 dB with respect to an input signal voltage of 200 μV. This might sound like only an average noise figure compared to that attainable with the moving magnet preamp, but it should be remembered that the noise generated by the cartridge itself is of this order of magnitude!

Another point worth mentioning here is that all the noise figures quoted so far in this article are flat or unweighted measurements. This means that the measurement was carried out with a noise and distortion analyser with a flat frequency response over the quoted noise bandwidth, usually 20 kHz. This is convenient and meaningful for the analysis of electrical circuits so long as the frequency distribution of the noise is also known. Probably the most useful way to quote noise figures at audio frequencies, however, is to graph noise circuits of nV/√Hz against frequency. The problem with flat noise measurements is that the human hearing mechanism does not detect all frequencies with equal sensitivity. For example, a noise generator with a high average noise voltage in the 100 Hz to 1 kHz band. To overcome this problem the frequency response of the measured equipment can be modified to accent or 'weight' the appropriate frequency bands. The most common weighting curve used in audio measurement is shown in Figure 6. The use of A-weighting gives a better indication of the apparent loudness of a noise voltage than do unweighted ('flat') measurements, and this is the reason almost all manufacturers quote A-weighted noise figures. A-weighted noise measurements for the Series 5000 MM and MC stages are quoted with specifications elsewhere in this article.



ETI-478MC PARTS LIST FOR STEREO PC BOARD

Resistors all 1/4W metal film, 5% unless noted otherwise	C7,C107 100p ceramic
R1,R101 220R	C8,C108 220p ceramic
R2,R102,R11,R111,R12,R112 . . . 100R	C9,C109 10u 16 V electrolytic
R3,R103 39R	C10,C110, C11,C111 47u 25 V electrolytic
R4,R104,R5,R105 . . . 390R	Semiconductors
R6,R106 22k	Q1,Q101,Q2,Q102 . BC549
R7,R107 6R8	Integrated circuits
R8,R108 680R	IC1, IC101 LM394CH
R9,R109 3k3	IC2, IC102 NE5534AN
R10,R110 1k	Miscellaneous
Capacitors	L1 Two turns on ferrite balun core, Neosid type 1050/2/F14 or 42-002-31
C1,C101,C3,C103 . . . 10n greencap	ETI-478MC pc board; shielded cable; assorted mounting hardware, etc.
C2,C102,C5,C105 . . . 1n greencap	
C4,C104 470u 16 V electrolytic	
C6,C106 680p ceramic	

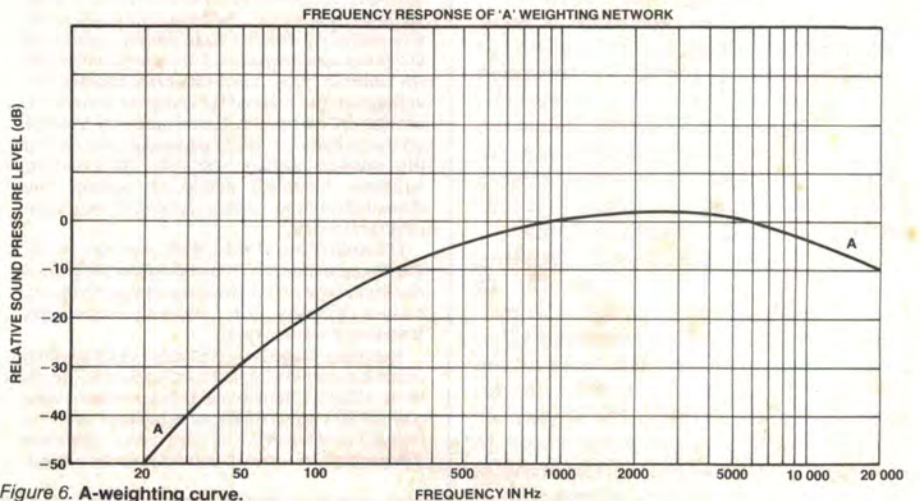


Figure 6. A-weighting curve.

